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A polyconvex extension of the logarithmic Hencky strain energy

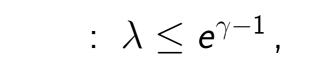
Robert J. Martin and Ionel-Dumitrel Ghiba and Patrizio Neff

1 Non-ellipticity of the Hencky energy

It has been known for a while that the *quadratic* isotropic Hencky energy W_H with [6]

where

 $\int -(\gamma-1)^2$



3 Valanis-Landel energies

We can apply the same extension method to the more general case of *Valanis-Landel type* energy functions, i.e. to functions of the form

$$W_{
m H}(F) = \mu \, \| {
m dev}_n \log U \|^2 + rac{\kappa}{2} \, [{
m tr}(\log U)]^2 \,, \qquad (1)$$

which is based on the logarithmic strain measures [9, 5]

 $||\text{dev}_n \log U||^2$, $||\log U||^2$ and $[\text{tr}(\log U)]^2$, (2) where $U = \sqrt{F^T F}$ denotes the Biot stretch tensor corresponding to the deformation gradient F, is not overall rank-one convex [8, 3]. Rank-one convexity (or Legendre-Hadamard ellipticity) is a necessary condition for polyconvexity, which, in turn, is an essential requirement for the applicability of existence proofs based on the direct methods of the calculus of variations [1]. This shortcoming raises some concern regarding the suitability of the Hencky model in finite element methods, although (rather large) ellipticity domains of the Hencky energy have been determined explicitly [3, 4].

Moreover [7], for $n \ge 3$, there exists no strictly monotone function $\Psi: [0, \infty) \to \mathbb{R}$ such that either of the energy functions $W: \operatorname{GL}^+(n) \to \mathbb{R}$ with

 $W(F) = \Psi(\|\log U\|^2)$ or $W(F) = \Psi(\|\operatorname{dev}_n \log U\|^2)$

is elliptic. If Ψ is additionally twice-differentiable, then there exists no smooth function W_{vol} : $(0, \infty) \to \mathbb{R}$ such that the energy W: $\text{GL}^+(n) \to \mathbb{R}$ with

$$arphi_{\gamma}(\lambda) = \left\{ egin{array}{ll} \ln^2(\lambda) + (2 - 2\gamma) \ln(\lambda) & : e^{\gamma - 1} < \lambda < e^{\gamma} \ -\gamma^2 + 2\gamma + rac{2}{e^{\gamma}}(e^{\lambda - e^{\gamma}} - 1) & : e^{\gamma} \leq \lambda \,. \end{array}
ight.$$

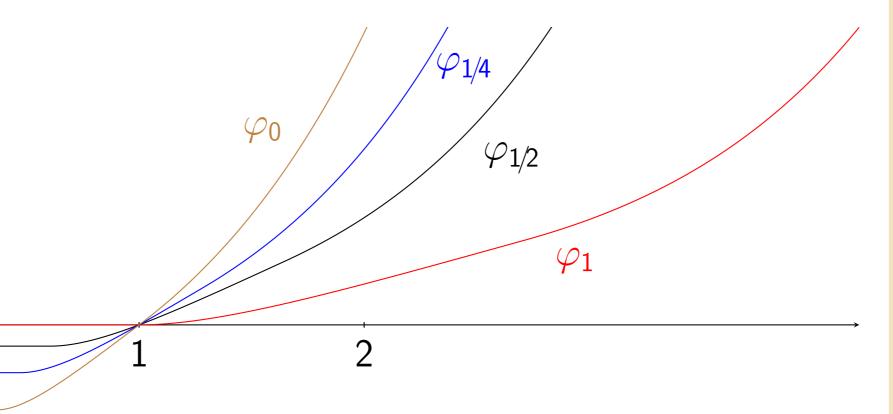


Figure 1: The function φ_{γ} for different values of γ .

Proposition. Let W_H denote the quadratic Hencky energy, given by

$$W_H(F) = \mu \|\log \sqrt{F^T F}\|^2 + \frac{\Lambda}{2} [\operatorname{tr}(\log \sqrt{F^T F})]^2,$$

where μ is the shear modulus and Λ is the first Lamé parameter. If $\Lambda \geq 0$, then the restriction of W_H to the set

$$W_{ extsf{VL}}\colon\operatorname{GL}^+(n) o\mathbb{R}$$
 , $W_{ extsf{VL}}(F)=\sum_{i=1}^n w(\lambda_i)$

with a scalar function $w: (0, \infty) \to \mathbb{R}$. Functions of this type were suggested by Valanis and Landel [11] as a general hyperelastic model for *incompressible* materials, but are often coupled additively with volumetric energy terms in order to obtain elastic models for compressible materials (including the quadratic Hencky energy $W_{\rm H}$ as well as Ogden's classical material model [10]). Note that the energy $W_{\rm VL}$ can only be compatible with linear elasticity at the identity 1 if w(1) = 0, w'(1) = 0 and w''(1) > 0; the latter two conditions represent the requirements of a stress-free reference configuration and ellipticity at 1, respectively.

Proposition. Let $w \in C^2((0, \infty))$ such that w'(1) = 0 and w''(1) > 0. Then the function

$$W_{ extsf{VL}}\colon\operatorname{GL}^+(n) o\mathbb{R}\,,\quad W_{ extsf{VL}}(F)=\sum_{i=1}^nw(\lambda_i)$$

has a polyconvex extension from a neighborhood of the

 $W(F) = \Psi(\|\operatorname{dev}_n \log U\|^2) + W_{\operatorname{vol}}(\det F)$

is Legendre-Hadamard elliptic.

2 A polyconvex extension of the classical Hencky energy

In order to find an elliptic energy function which approximates (or, better yet, is identical to) the Hencky strain energy in the small-strain range, we adapt an approach by Ball, Muite, Schryvers and Tirry [2] to construct a polyconvex (and thus rank-one convex) extension of the quadratic-logarithmic Hencky energy (1) and, more generally, for suitable energy expressions of the Valanis-Landel type. In addition, the extension of the Hencky energy considered here is (unconditionally) coercive, which implies an immediate applicability of the direct methods of the calculus of variations to prove the existence of energy minimizers under appropriate boundary conditions.

Lemma. For $\gamma \leq 1$, let

 $\mathcal{S}_{\gamma} := \{F \in \mathsf{GL}^+(n) \mid e^{\gamma-1} < \lambda < e^{\gamma} \text{ for each singular value } \lambda \text{ of } F\}.$

 $\mathcal{S}_{1/3} = \{F \in \mathsf{GL}^+(n) \mid e^{-2/3} < \lambda < e^{1/3} ext{ for each singular value } \lambda ext{ of } F\}$

has a polyconvex extension to $GL^+(n)$.

The energy \widetilde{W}_{H} is polyconvex as well as coercive and bounded below, allowing for a direct application of Ball's classical result on the existence of minimizers [1].

Proposition. Let $\Omega \subset \mathbb{R}^n$ be a bounded smooth domain, Γ_D be a non-empty and relatively open part of the boundary $\partial\Omega$ and $\varphi_0 \in W^{1,q}(\Omega)$ for some q > 1 such that $\int_{\Omega} \widetilde{W}_{H}(\nabla \varphi_0(x)) dx < \infty$. Then the minimization problem

$$\min_{\substack{\varphi \in W^{1,p}(\Omega) \\ \varphi \mid_{\Gamma_D} = \varphi_0}} \int_{\Omega} \widetilde{W}_{\mathsf{H}}(\nabla \varphi(x)) \, \mathsf{d}x$$

admits at least one solution $\widehat{\varphi} \in W^{1,p}(\Omega)$.

identity F = 1 to $GL^+(n)$.

4 **References**

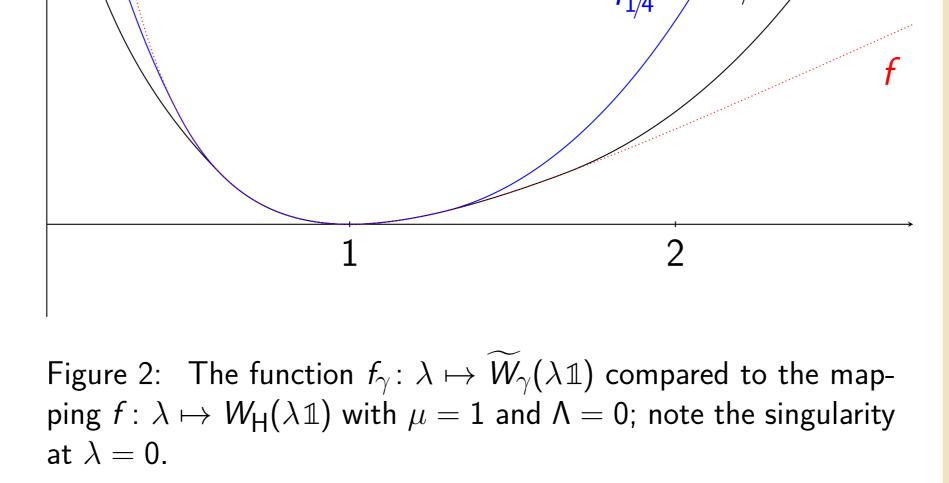
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$$W \colon \mathcal{S}_{\gamma} o \mathbb{R}$$
, $W(F) = \|\log \sqrt{F^{\intercal}F}\|^2 = \sum_{i=1}^n \ln^2(\lambda_i)$

has a polyconvex extension \widetilde{W}_{γ} : $\operatorname{GL}^+(n) \to \mathbb{R}$ to $\operatorname{GL}^+(n)$, which is given by

$$\widetilde{W}_{\gamma}(F) = \sum_{i=1}^{n} \varphi_{\gamma}(\lambda_i) - (2 - 2\gamma) \ln(\det F),$$



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