

# **Comparison of Models for Finite Plasticity**

**A numerical study**

**Patrizio Neff**

**and**

**Christian Wieners**

**California Institute of Technology**

**(Universität Darmstadt)**

**Universität Augsburg**

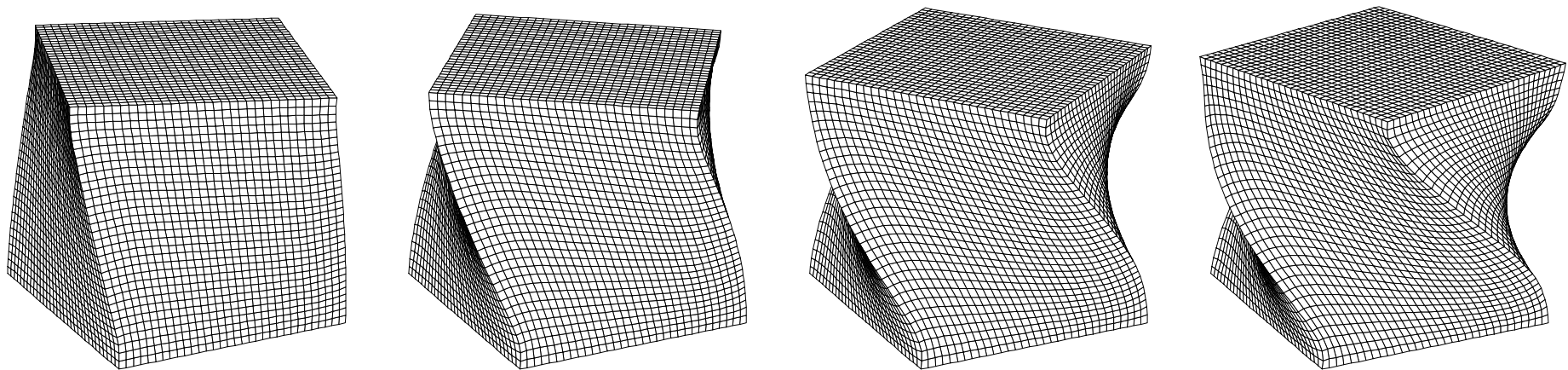
**(Universität Heidelberg)**

# Motivation

---

For a wide range of applications of elasto-plasticity theory to crystalline aggregates like metals it is known that the elastic region is very small implying that elastic strains remain small.

A frame indifferent approximative model for this situation has been developed where the remaining nonlinearity (in case of small elastic strains) in the balance equation is shifted into an appended evolution equation. It has been shown that the new model in a viscous setting is locally well posed.



Here, we compare *numerically* the behavior of the new model with other well known models of finite plasticity which are based on different elastic free energies.

# Models

## Small elastic strains

---

We assume that the deformation tensor  $F$  has a multiplicative decomposition  $F = F_e F_p$ ; we set  $P = F_p^{-1}$  and  $F_e = FP$ . We consider the polar decomposition of the elastic part

$$F_e = \text{polar}(F_e) U_e \quad \text{with} \quad \text{polar}(F_e) \in SO(3) \quad \text{and} \quad U_e = \sqrt{F_e^T F_e}.$$

For small elastic strains (i. e.  $|F_e^T F_e - \text{id}|$  is small), also  $|F_e - \text{polar}(F_e)|$  is small, and

$$F_e^T \text{polar}(F_e) + \text{polar}(F_e)^T F_e - 2\text{id} \left( = 2(U_e - \text{id}) \right) \quad \text{is close to} \quad F_e^T F_e - \text{id}.$$

In the new model,  $\text{polar}(F_e)$  is approximated by  $R \in SO(3)$  satisfying

$$\dot{R} = \nu(|\text{skew}(F_e R^T)|) \text{skew}(F_e R^T) R.$$

**Lemma.** For  $R_0 \in SO(3)$  with  $|R_0 - \text{polar}(F_e)| < 8$ , the evolution equation

$$\dot{R} = \text{skew}(F_e R^T) R, \quad R(0) = R_0$$

has a global solution  $R(t)$  in  $SO(3)$  satisfying  $\lim_{t \rightarrow \infty} R(t) = \text{polar}(F_e)$ .

# Material models

---

**M1** The new model (with appended rotations  $R$ )

**M2** The exact polar model (using  $R = \text{polar}(F_e)$ )

**M3** St. Venant–Kirchhoff model (straight forward finite strain extension)

**M4** A Neo-Hooke model (reference model)

**M5** A Hencky model with logarithmic strains (fits to exponential increment)

For comparison, we also consider a classical model with linearized plastic strain.

**M6** Perfect plasticity (small strain limit)

**M7** Viscoplasticity (measuring the viscous effect)

# Internal variables, strain and energy

|                                   | internal variables $Z$ | strain $C = \text{id} + 2E$                                     | energy $W$                                                                                  |
|-----------------------------------|------------------------|-----------------------------------------------------------------|---------------------------------------------------------------------------------------------|
| <b>M1</b><br>new model            | $(P, R)$               | $F_e^T R + R^T F_e - \text{id}$                                 | $\mu  E ^2 + \frac{\lambda}{2} \text{trace}(E)^2$                                           |
| <b>M2</b><br>exact polar model    | $P$                    | $F_e^T \text{polar}(F_e) + \text{polar}(F_e)^T F_e - \text{id}$ | $\mu  E ^2 + \frac{\lambda}{2} \text{trace}(E)^2$                                           |
| <b>M3</b><br>St. Venant–Kirchhoff | $P$                    | $F_e^T F_e$                                                     | $\mu  E ^2 + \frac{\lambda}{2} \text{trace}(E)^2$                                           |
| <b>M4</b><br>Neo-Hooke            | $P$                    | $F_e^T F_e$                                                     | $\frac{\mu}{2}  F_e ^2 + \frac{\lambda}{2} \det(C) - \frac{2\mu + \lambda}{2} \log \det(C)$ |
| <b>M5</b><br>Hencky               | $P$                    | $F_e^T F_e$                                                     | $\frac{\mu}{4}  \text{dev}(\ln C) ^2 + \frac{\kappa}{8} (\text{trace}(\ln C))^2$            |
| <b>M6</b><br>Perfect plasticity   | $\varepsilon^p$        | $F + F^T - \text{id} - 2\varepsilon^p$                          | $\mu  E ^2 + \frac{\lambda}{2} \text{trace}(E)^2$                                           |
| <b>M7</b><br>Viscoplasticity      | $\varepsilon^p$        | $F + F^T - \text{id} - 2\varepsilon^p$                          | $\mu  E ^2 + \frac{\lambda}{2} \text{trace}(E)^2$                                           |

# Finite strain models

---

For all models, we set

$$\begin{aligned} S &= D_F W(F, Z) \\ \Sigma &= F_e^T S P^{-T}, \\ D &= \text{dev}(\text{sym}(\Sigma)), \\ P^{-1} \dot{P} &= -\gamma(|D|) \frac{D}{|D|}, \\ \gamma(d) &\in \partial \hat{\chi}_r(d). \end{aligned}$$

For **M2**, **M3**, **M4**, **M5**, **M6** we use  $\hat{\chi}_\infty(d) = \begin{cases} 0 & d \leq K_0, \\ \infty & d > K_0. \end{cases}$

In **M1** and **M7** this is relaxed to  $\hat{\chi}_r(d) = \frac{1}{\eta} \frac{1}{r+1} \max\{0, d - K_0\}^{r+1}, \quad r \in [1, \infty).$

## Quasi-static models

---

Let  $\Omega \subset \mathbf{R}^3$  be the reference domain, let  $\Gamma_D \subset \partial\Omega$  be the Dirichlet boundary part (with non-vanishing measure), let  $\Gamma_N = \partial\Omega \setminus \Gamma_D$  be the Neumann boundary part, let  $n$  be the outer normal on  $\Gamma_N$ , and let  $[0, T]$  be a time interval.

The constitutive equations are complemented by

$$\begin{aligned} -\operatorname{div} S &= f && \text{in } [0, T] \times \Omega, \\ S \cdot n &= g && \text{on } [0, T] \times \Gamma_N, \\ F &= \operatorname{id} + \nabla u && \text{in } [0, T] \times \Omega, \\ u &= d && \text{on } [0, T] \times \Gamma_D \end{aligned}$$

for prescribed volume forces  $f$ , surface loads  $g$ , and Dirichlet boundary displacements  $d$ .

**Theorem (Neff).** In the pure Dirichlet case, for smooth data and  $r > 6$ , a time  $T > 0$  exists such that the new model **M1** has a unique solution

$$(u, P, R) \in C([0, T], H^{2,5}(\Omega, \mathbf{R}^3)) \times C^1([0, T], H^{4,2}(\Omega, SL(3)), H^{4,2}(\Omega, SO(3))).$$



## Variational problem

---

The material is completely determined by the evolution of the internal variables

$$\dot{Z} \in H(F, Z)$$

(depending on a history function  $H$ ) and the stress response

$$S = S(F, Z).$$

The displacement  $u$  is determined by the boundary values  $u = d$  on  $\Gamma_D$  and

$$\int_{\Omega} S(\text{id} + \nabla u, Z) \cdot \nabla v \, dx = \int_{\Omega} f \cdot v \, dx + \int_{\Gamma_N} g \cdot v \, ds$$

for all  $v$  with  $v = 0$  on  $\Gamma_D$ .

# **Discretization**

## Incremental problem

---

For a time series  $0 = t_0 < t_1 < t_2 < \dots < t_n = T$  with time increments  $\tau_n = t_n - t_{n-1}$ , the incremental problem is completely determined by a time discretization

$$Z_n = H_n(\tau_n, F_n, Z_{n-1})$$

of the evolution of the internal variables.

This defines the incremental stress response

$$S = S(F_n, H_n(\tau_n, F_n, Z_{n-1})).$$

The displacement  $u_n$  is determined by the boundary values  $u_n = d(t_n)$  on  $\Gamma_D$  and

$$\int_{\Omega} S(\text{id} + \nabla u_n, H_n(\tau_n, \text{id} + \nabla u_n, Z_{n-1})) \cdot \nabla v \, dx = \int_{\Omega} f \cdot v \, dx + \int_{\Gamma_N} g \cdot v \, ds$$

for all  $v$  with  $v = 0$  on  $\Gamma_D$ , depending on the history variables  $Z_{n-1}$ .

## Exponential update in finite plasticity

---

The evolution problem in  $t \in [t_{n-1}, t_n]$

$$\dot{P}(t) = -\gamma(t)P(t)\frac{D(t)}{|D(t)|} \quad \text{with} \quad \begin{aligned} P(t_{n-1}) &= P_{n-1} \\ D(t) &= \text{dev}(\text{sym}(P^T(t)F^T(t)S(t)P^{-T}(t))), \\ \gamma(t) &\in \partial\hat{\chi}_r(|D(t)|) \end{aligned}$$

is discretized by the linear evolution equation

$$\dot{P}(t) = -\gamma_n P(t)\frac{D_n}{|D_n|} \quad \text{with} \quad \begin{aligned} P(t_{n-1}) &= P_{n-1} \\ D_n &= \text{dev}(\text{sym}(P_n^T F_n^T S_n P_n^{-T})), \\ \gamma_n &\in \partial\hat{\chi}_r(|D_n|); \end{aligned}$$

this gives (depending on  $\gamma_n$  and  $D_n$ )

$$P_n = P_{n-1} \exp\left(-\tau_n \gamma_n \frac{D_n}{|D_n|}\right)$$

Note that this procedure preserves the constraint  $\det(P_n) = 1$ .

## History increment for M2

---

For given  $F_n$  and  $P_{n-1}$ , compute  $D_n^{\text{trial}} = D(F_n, P_{n-1})$ . We consider two cases:

for  $|D_n^{\text{trial}}| \leq K_0$  we set  $P_n = H_n(F_n, P_{n-1}) = P_{n-1}$ ;

for  $|D_n^{\text{trial}}| > K_0$  we compute  $(P_n, \gamma_n)$  such that

$$\begin{pmatrix} P_n - P_{n-1} \exp\left(-\tau_n \gamma_n \frac{D(F_n, P_n)}{|D(F_n, P_n)|}\right) \\ |D(F_n, P_n)| - K_0 \end{pmatrix} = 0$$

with

$$\begin{aligned} D(F, P) = & \operatorname{dev}\left(P^T F^T \operatorname{polar}(FP) (\mu(P^T F^T \operatorname{polar}(FP) + \operatorname{polar}(FP)^T F P - 2 \operatorname{id}) \right. \\ & \left. + \frac{\lambda}{2} \operatorname{trace}(P^T F^T \operatorname{polar}(FP) + \operatorname{polar}(FP)^T F P - 2 \operatorname{id}))\right). \end{aligned}$$

## History increment for M1

---

For given  $F_n$  and  $Z_{n-1} = (P_{n-1}, R_{n-1})$ , compute  $D_n^{\text{trial}} = D(F_n, P_{n-1}, R_{n-1})$ .

For  $|D_n^{\text{trial}}| > K_0$  we compute  $(P_n, R_n, \gamma_n, \mathbf{v}_n)$  such that

$$\begin{pmatrix} P_n - P_{n-1} \exp\left(-\tau_n \gamma_n \frac{D(F_n, P_n, R_n)}{|D(F_n, P_n, R_n)|}\right) \\ R_n - \exp\left(\tau_n \mathbf{v}_n \text{skew}(F_n P_n R_n^T)\right) R_{n-1} \\ \gamma_n - \frac{1}{\eta} (|D(F_n, P_n, R_n)| - K_0)^r \\ \mathbf{v}_n - \frac{1}{\eta} (1 + |\text{skew}(F_n P_n R_n^T)|)^r \end{pmatrix} = 0$$

with

$$D(F, P, R) = \text{dev}\left(\left(\text{id} + \frac{1}{2}(P^T F^T R + R^T F P - 2\text{id})\right)\right. \\ \left. \left(\mu(P^T F^T R + R^T F P - 2\text{id}) + \frac{\lambda}{2} \text{trace}(P^T F^T R + R^T F P - 2\text{id})\right)\right).$$

## Finite element setting

---

For finite elements  $V_h$  with nodal points  $P_h$  set  $V_h(d) = \{v \mid v(P) = d(P), P \in P_h \cap \Gamma_D\}$ .  
For given material history  $Z_{n-1}$ , find  $u_n \in V_h(d_n)$  such that

$$\int_{\Omega} S(\text{id} + \nabla u_n, H_n(\tau_n, \text{id} + \nabla u_n, Z_{n-1})) \cdot \nabla v \, dx = \int_{\Omega} f \cdot v \, dx + \int_{\Gamma_N} g \cdot v \, ds, \quad v \in V_h(0).$$

This is solved by a Newton method: for given  $u_n^m \in V_h(d_n)$  find  $w \in V_h(0)$  such that

$$\int_{\Omega} C_n^m \nabla w \cdot \nabla v \, dx = \int_{\Omega} f \cdot v \, dx + \int_{\Gamma_N} g \cdot v \, ds - \int_{\Omega} S(\text{id} + \nabla u_n^m, Z(\tau_n, \text{id} + \nabla u_n^m, Z_{n-1})) \cdot \nabla v \, dx$$

for  $v \in V_h(0)$  with the approximated consistent tangent operator  $C_n^m \approx D_F S = D_F^2 W$

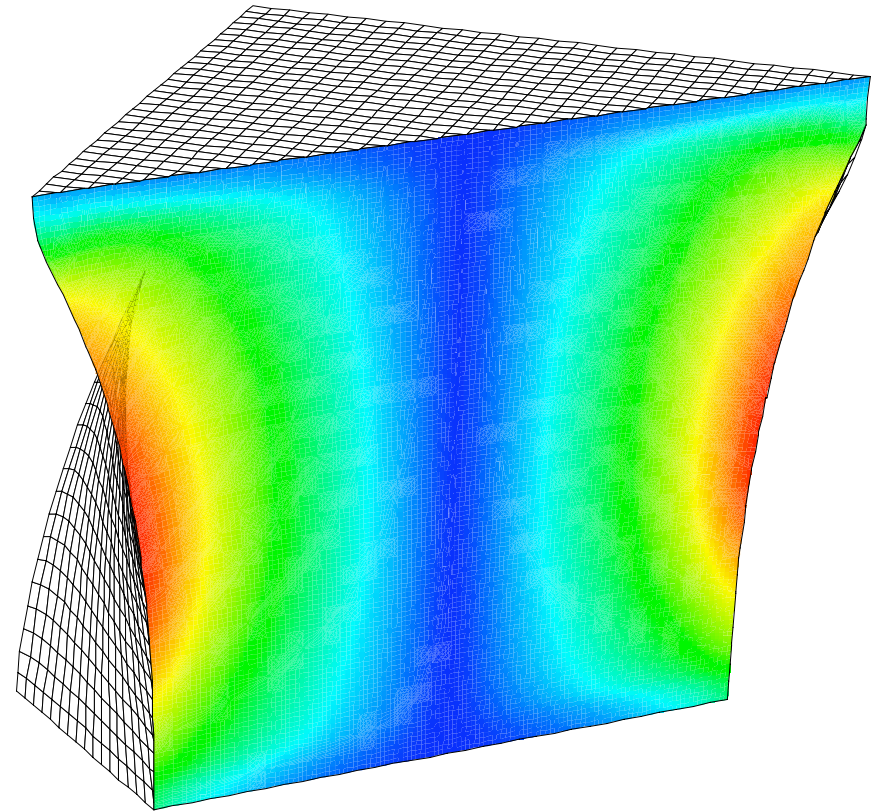
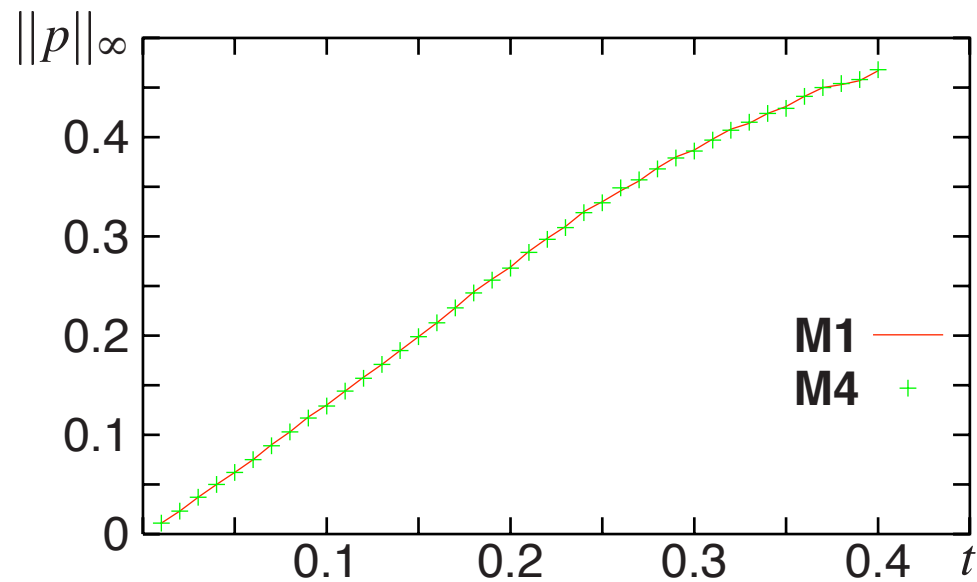
$$\begin{aligned} C_n^m \nabla w &= \frac{1}{\delta} \left( S(\text{id} + \nabla u_n^m + \delta \nabla w, Z(\tau_n, \text{id} + \nabla u_n^m + \delta \nabla w, Z_{n-1})) \right. \\ &\quad \left. - S(\text{id} + \nabla u_n^m, Z(\tau_n, \text{id} + \nabla u_n^m, Z_{n-1})) \right) \quad (\delta > 0). \end{aligned}$$

# **Numerical Comparison**



# Torsion test

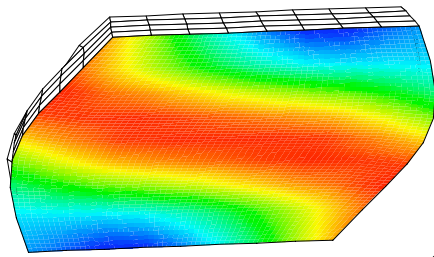
---



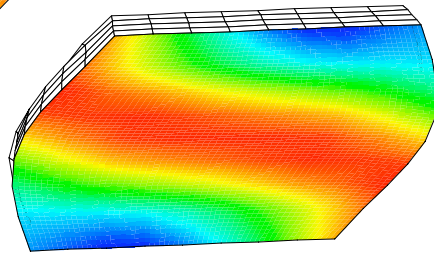
Evolution the equivalent plastic strain  $p(t) = \sqrt{\frac{2}{3}} \int_0^t \gamma ds$  for **M1** and **M4**

# Shear and compression test

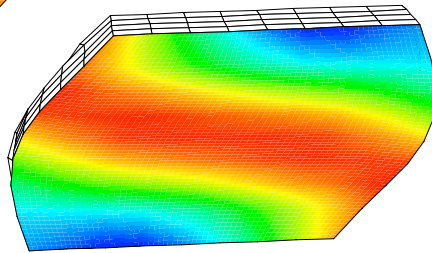
---



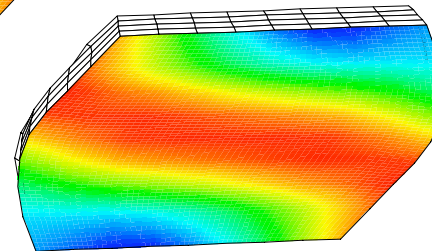
**M2**



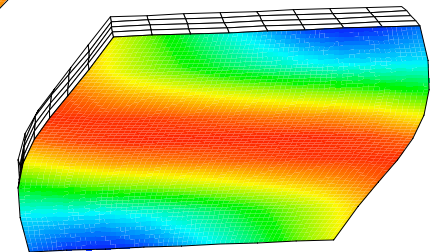
**M3**



**M4**



**M5**



**M6**

Distribution of the equivalent plastic strain  
for a deformation of 25%

## Shear and compression test

| $\ p(t)\ _{\infty, \Omega}$    | <b>M2</b> | <b>M3</b> | <b>M4</b> | <b>M5</b> | <b>M6</b> |
|--------------------------------|-----------|-----------|-----------|-----------|-----------|
| $\ p(0.02)\ _{\infty, \Omega}$ | 0.060     | 0.060     | 0.060     | 0.060     | 0.053     |
| $\ p(0.1)\ _{\infty, \Omega}$  | 0.289     | 0.289     | 0.289     | 0.289     | 0.251     |
| $\ p(0.3)\ _{\infty, \Omega}$  | 0.878     |           | 0.877     | 0.879     | 0.662     |

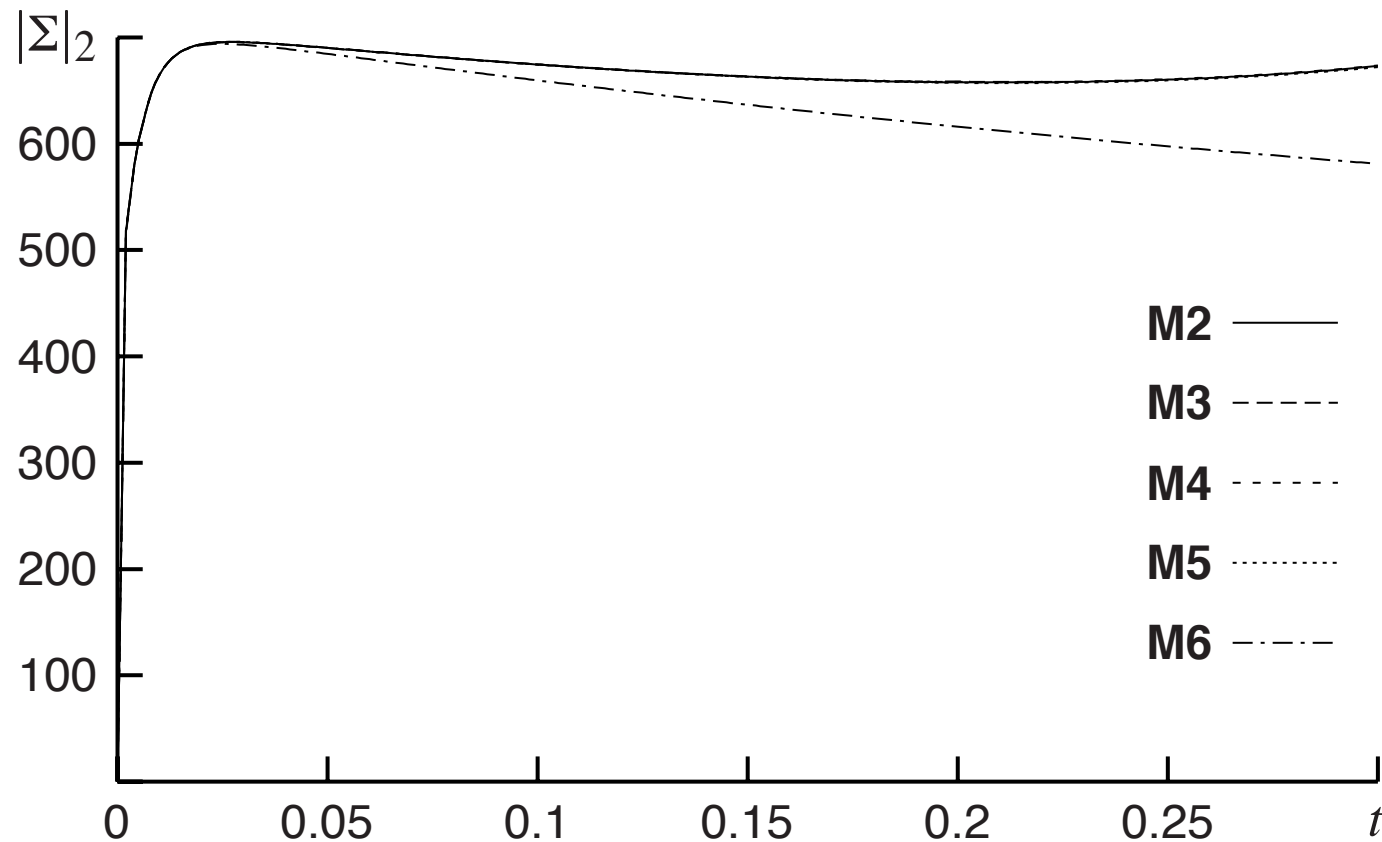
Results for the finite strain models **M2**, **M3**, **M4**, **M5** (and for comparison perfect plasticity **M6**) for deformations of 2%, 10% and 30% (for **M3** and for  $t > 0.2$  the numerical algorithm is converged only on coarser meshes).

As a cross check for nonlinear elasticity, we present the stress invariants for a deformation of 0.1%; the stress response is purely elastic in this case (where  $F_p \equiv 0$ ).

|                     | <b>M2</b> | <b>M3</b> | <b>M4</b> | <b>M5</b> | <b>M6</b> |
|---------------------|-----------|-----------|-----------|-----------|-----------|
| $\Sigma_I(x_m)$     | 15.7      | 15.8      | 15.7      | 15.7      | 15.6      |
| $\Sigma_{II}(x_m)$  | -7.8      | -7.7      | -7.8      | -7.8      | -7.8      |
| $\Sigma_{III}(x_m)$ | -263.6    | -263.1    | -263.9    | -264.2    | -264.0    |

## Shear and compression test

---



The evolution of the maximal stress  $|\Sigma|_2 = \max\{\Sigma_I(x_m), \Sigma_{II}(x_m), \Sigma_{III}(x_m)\}$  for the models **M2**, **M3**, **M4**, **M5**, and **M6**. Here,  $x_m = (0.5, 0.5, 0.5)$  denotes the midpoint of  $\Omega$ , and  $\Sigma_I$ ,  $\Sigma_{II}$ ,  $\Sigma_{III}$  are the eigenvalues of the stress tensor  $\Sigma$  (resp.  $\text{sym}(\Sigma)$  for **M2**).

## Shear and compression test - the effect of viscoplastic relaxation

|                                          | <b>M6</b> |        | <b>M7</b> |        |
|------------------------------------------|-----------|--------|-----------|--------|
| $\ p\ _{\infty,\Omega}$                  | 0.027     | 0.057  | 0.027     | 0.057  |
| $\ \text{dev}(\sigma)\ _{\infty,\Omega}$ | 450       | 450    | 450.5     | 450.5  |
| $\Sigma_I(x_m)$                          | -665.2    | -692.7 | -665.8    | -693.3 |
| $\Sigma_{II}(x_m)$                       | -175.4    | -185.3 | -175.4    | -185.4 |
| $\Sigma_{III}(x_m)$                      | -68.5     | -106.3 | -68.4     | -106.3 |

Results for the perfect plasticity model **M6** and the viscoplastic regularization **M7** for a deformation of 1% and 2%.

|                                          | <b>M2</b> |        | <b>M1</b> |        |
|------------------------------------------|-----------|--------|-----------|--------|
| $\ p\ _{\infty,\Omega}$                  | 0.289     | 0.878  | 0.289     | 0.875  |
| $\ \text{dev}(\Sigma)\ _{\infty,\Omega}$ | 450       | 450    | 450.5     | 450.5  |
| $\Sigma_I(x_m)$                          | -674.6    | -673.1 | -682.8    | -697.4 |
| $\Sigma_{II}(x_m)$                       | -151.3    | -141.8 | -158.4    | -166.3 |
| $\Sigma_{III}(x_m)$                      | -99.2     | -104.1 | -107.3    | -127.0 |

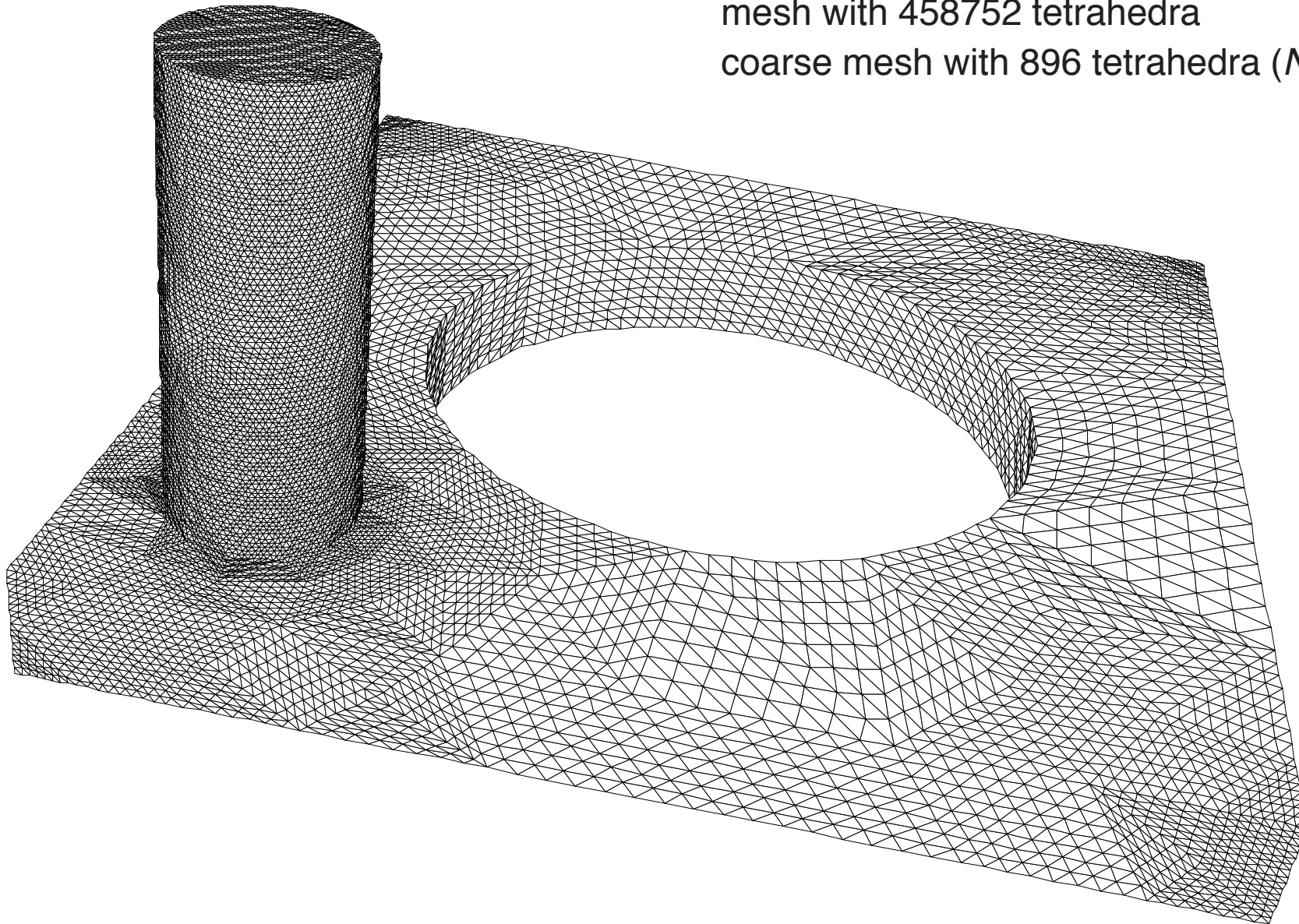
Results for the model **M2** and the viscoplastic regularization **M1** ( $t = 0.1$  and  $0.3$ ).

## Advanced application

---

mesh with 458752 tetrahedra

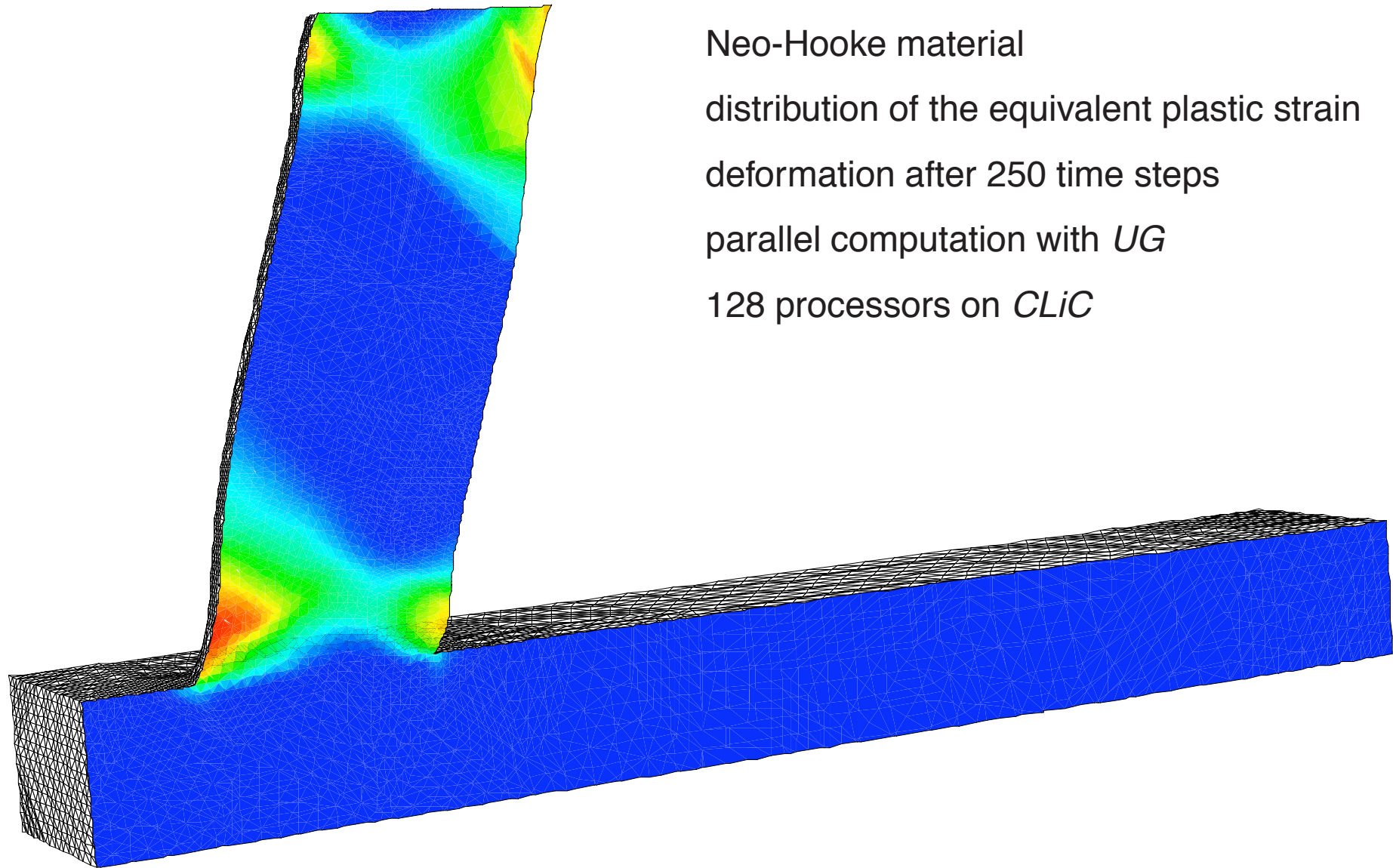
coarse mesh with 896 tetrahedra (*NETGEN*)





## Advanced application

---



Neo-Hooke material

distribution of the equivalent plastic strain

deformation after 250 time steps

parallel computation with *UG*

128 processors on *CLiC*