

Numerical Analysis for Maxwell Obstacle Problems in Faraday Shielding

jointly with Irwin Yousept

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Obstacle problem in Faraday shielding

Faraday shielding (recap): Effect of redirecting or blocking certain electric fields.

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Measurement of an electric field by an EMF-meter with and without Faraday shielding

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$$\mathbf{E}(t) \in K := \{\mathbf{v} \in L^2(\Omega) \mid |\mathbf{v}(x)| \leq d \text{ for a.e. } x \in \omega\}$$

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- $\Omega \subset \mathbb{R}^3$ and $\omega \subset\subset \Omega$ bounded, polyhedral Lipschitz domains
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- $(\mathbf{E}_0, \mathbf{H}_0) \in (\mathbf{K} \cap \mathbf{H}_0(\mathbf{curl})) \times \mathbf{H}(\mathbf{curl})$ initial data

Theorem

The obstacle problem (P) admits a unique solution

$$(\mathbf{E}, \mathbf{H}) \in W^{1,\infty}((0, T), L^2(\Omega) \times L^2(\Omega)) \cap L^\infty((0, T), \mathbf{H}_0(\mathbf{curl}) \times L^2(\Omega))$$

satisfying the local magnetic regularity

$$\mathbf{H}|_{\Omega \setminus \bar{\omega}} \in L^\infty((0, T), \mathbf{H}(\mathbf{curl}, \Omega \setminus \bar{\omega})).$$

↪ Result by Irwin Yousept ¹

¹I. Yousept. Well-posedness theory for electromagnetic obstacle problems. *J. Differential Equations*, 269(10):8855–8881, 2020

First Attempt: Mixed FEM and implicit Euler

- Denote by $\{\mathcal{T}_h\}_{h>0}$ a quasi-uniform family of triangulations, s.t.

$$\bar{\Omega} = \bigcup_{T \in \mathcal{T}_h} T, \quad \bar{\omega} = \bigcup_{T \in \mathcal{T}_h^\omega} T,$$

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- Time partition:

$$\tau = \frac{T}{N}, \quad t_n = n\tau \quad \forall n \in \{0, \dots, N\}, \quad N \in \mathbb{N}$$

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$$\left\{ \begin{array}{l} \text{Find } \{(E_h^n, H_h^n)\}_{n=1}^N \subset (K \cap \mathbf{ND}_h) \times \mathbf{DG}_h, \text{ s.t.} \\ \int_{\Omega} (\epsilon + \tau\sigma) E_h^n \cdot (v_h - E_h^n) + \tau^2 \mu^{-1} \text{curl } E_h^n \cdot \text{curl}(v_h - E_h^n) \, dx \\ \geq \int_{\Omega} (\tau f^n + E_h^{n-1}) \cdot (v_h - E_h^n) + \tau H_h^{n-1} \cdot \text{curl}(v_h - E_h^n) \quad \forall v_h \in K \cap \mathbf{ND}_h \\ H_h^n = H_h^{n-1} - \tau \mu^{-1} \text{curl } E_h^n. \end{array} \right.$$

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Drawback: High computational cost due to the requirement of a nonsmooth solver!

Second Attempt: Mixed FEM and Yee stepping

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 - the Amperé-Maxwell VI in (P) at $t_{n-\frac{1}{2}} := t_n - \frac{\tau}{2}$
 - the Faraday equation in (P) at t_n

²K. Yee. Numerical solution of initial boundary value problems involving Maxwell's equations in isotropic media. *IEEE Transactions on Antennas and Propagation*, 14(3):302–307, 1966

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²K. Yee. Numerical solution of initial boundary value problems involving Maxwell's equations in isotropic media. *IEEE Transactions on Antennas and Propagation*, 14(3):302–307, 1966

$$(P_{N,h}) \quad \left\{ \begin{array}{l} \text{Find } \{(E_h^{n-\frac{1}{2}}, H_h^{n+\frac{1}{2}})\}_{n=1}^N \subset (K \cap \mathbf{DG}_h) \times \mathbf{ND}_h \text{ s.t.} \\ \int_{\Omega} \epsilon \delta E_h^n \cdot (\mathbf{v}_h - E_h^{n-\frac{1}{2}}) + \sigma E_h^{n-\frac{1}{2}} \cdot (\mathbf{v}_h - E_h^{n-\frac{1}{2}}) - \text{curl } H_h^{n-\frac{1}{2}} \cdot (\mathbf{v}_h - E_h^{n-\frac{1}{2}}) dx \\ \geq \int_{\Omega} \mathbf{f}_h^{n-\frac{1}{2}} \cdot (\mathbf{v}_h - E_h^{n-\frac{1}{2}}) dx \quad \forall \mathbf{v}_h \in K \cap \mathbf{DG}_h \quad \forall n \in \{1, \dots, N\} \\ \int_{\Omega} \mu \delta H_h^{n+\frac{1}{2}} \cdot \mathbf{w}_h + E_h^n \cdot \text{curl } \mathbf{w}_h dx = 0 \quad \forall \mathbf{w}_h \in \mathbf{ND}_h \quad \forall n \in \{1, \dots, N\}, \end{array} \right.$$

with

$$\delta E_h^n := \frac{E_h^n - E_h^{n-1}}{\tau}, \quad \delta H_h^{n+\frac{1}{2}} := \frac{H_h^{n+\frac{1}{2}} - H_h^{n-\frac{1}{2}}}{\tau}, \quad E_h^n := 2E_h^{n-\frac{1}{2}} - E_h^{n-1}.$$

$$(P_{N,h}) \quad \left\{ \begin{array}{l} \text{Find } \{(\mathbf{E}_h^{n-\frac{1}{2}}, \mathbf{H}_h^{n+\frac{1}{2}})\}_{n=1}^N \subset (\mathbf{K} \cap \mathbf{DG}_h) \times \mathbf{ND}_h \text{ s.t.} \\ \int_{\Omega} \epsilon \delta \mathbf{E}_h^n \cdot (\mathbf{v}_h - \mathbf{E}_h^{n-\frac{1}{2}}) + \sigma \mathbf{E}_h^{n-\frac{1}{2}} \cdot (\mathbf{v}_h - \mathbf{E}_h^{n-\frac{1}{2}}) - \text{curl } \mathbf{H}_h^{n-\frac{1}{2}} \cdot (\mathbf{v}_h - \mathbf{E}_h^{n-\frac{1}{2}}) dx \\ \geq \int_{\Omega} \mathbf{f}_h^{n-\frac{1}{2}} \cdot (\mathbf{v}_h - \mathbf{E}_h^{n-\frac{1}{2}}) dx \quad \forall \mathbf{v}_h \in \mathbf{K} \cap \mathbf{DG}_h \quad \forall n \in \{1, \dots, N\} \\ \int_{\Omega} \mu \delta \mathbf{H}_h^{n+\frac{1}{2}} \cdot \mathbf{w}_h + \mathbf{E}_h^n \cdot \text{curl } \mathbf{w}_h dx = 0 \quad \forall \mathbf{w}_h \in \mathbf{ND}_h \quad \forall n \in \{1, \dots, N\}, \end{array} \right.$$

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Note: Obstacle discretization at $t_{n-\frac{1}{2}}$ rather than t_n & L^2 -structure

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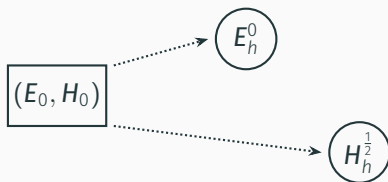
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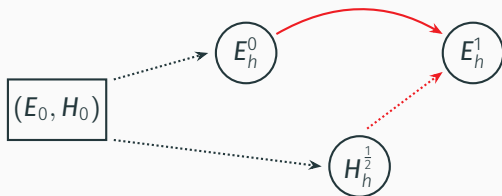
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(E_0, H_0)

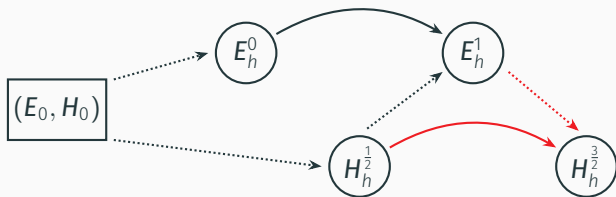
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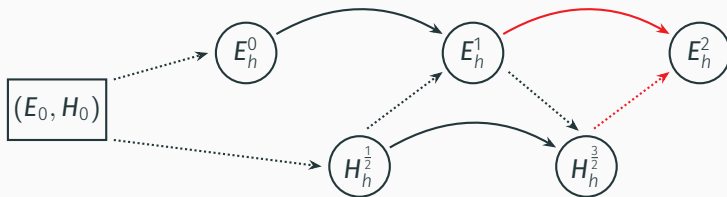
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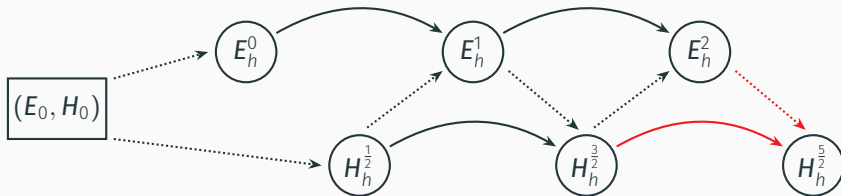
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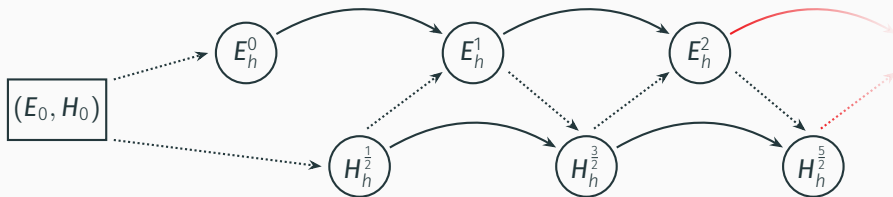
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Theorem

The problem $(P_{N,h})$ admits a unique solution $\{(E_h^{n-\frac{1}{2}}, H_h^{n+\frac{1}{2}})\}_{n=1}^N \subset (K \cap \mathbf{DG}_h) \times \mathbf{ND}_h$ with

$$E_h^{n-\frac{1}{2}} = \begin{cases} \frac{dg_h^{n-\frac{1}{2}}}{|g_h^{n-\frac{1}{2}}|} & \text{on } \mathcal{M}_h^{n-\frac{1}{2}} \\ \left(\frac{2\epsilon}{\tau} + \sigma\right)^{-1} g_h^{n-\frac{1}{2}} & \text{on } \Omega \setminus \mathcal{M}_h^{n-\frac{1}{2}}, \end{cases}$$

with right-hand sides and strict superlevel sets

$$g_h^{n-\frac{1}{2}} := f_h^{n-\frac{1}{2}} + \operatorname{curl} H_h^{n-\frac{1}{2}} + \frac{2\epsilon}{\tau} E_h^{n-1} \quad \text{and} \quad \mathcal{M}_h^{n-\frac{1}{2}} := \left\{ x \in \omega \mid \left(\frac{2\epsilon}{\tau} + \sigma\right)^{-1} |g_h^{n-\frac{1}{2}}(x)| > d \right\}.$$

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We recall the inverse estimate

$$\exists C_{\text{inv}} > 0 \text{ s.t. } \|\mathbf{curl} \mathbf{v}\|_{L^2(\Omega)} \leq \frac{C_{\text{inv}}}{h} \|\mathbf{v}\|_{L^2(\Omega)} \quad \forall \mathbf{v} \in \mathbf{ND}_h$$

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- additional regularity on the initial electric field

$$\mathbf{E}_0 \in \mathbf{K} \cap \mathbf{H}_0(\mathbf{curl}) \cap H^1(\Omega)$$

↪ Main ingredients for stability.

Theorem

There exists a constant $C > 0$ such that for every $N \in \mathbb{N}$ with $N \geq 2$ and $h > 0$ the unique solution to $(P_{N,h})$ satisfies

$$\max_{n \in \{1, \dots, N\}} \|\delta E_h^n\|_{L^2(\Omega)} + \max_{n \in \{2, \dots, N\}} \|\delta H_h^{n-\frac{1}{2}}\|_{L^2(\Omega)} \leq C$$

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Lack of global L^2 -stability for $\mathbf{curl} H_h^{n-\frac{1}{2}}$: Justified by low regularity issue in (P).

We set up the following **piecewise linear interpolations**

$$E_{N,h}: [0, T] \rightarrow \mathbf{DG}_h, \quad t \mapsto \begin{cases} E_h^0 & \text{if } t = 0 \\ E_h^{n-1} + (t - t_{n-1})\delta E_h^n & \text{if } t \in (t_{n-1}, t_n], \end{cases}$$

$$H_{N,h}: [0, T] \rightarrow \mathbf{ND}_h, \quad t \mapsto \begin{cases} H_h^{\frac{1}{2}} & \text{if } t = 0 \\ H_h^{n-\frac{1}{2}} + (t - t_{n-1})\delta H_h^{n+\frac{1}{2}} & \text{if } t \in (t_{n-1}, t_n] \text{ for } n \in \{1, \dots, N-1\} \\ H_h^{N-\frac{3}{2}} & \text{if } t \in (t_{n-1}, t_n] \text{ for } n = N. \end{cases}$$

Theorem

Under the stated CFL-condition, it holds that

$$\begin{aligned} (\mathbf{E}_{N,h}, \mathbf{H}_{N,h}) &\overset{*}{\rightharpoonup} (\mathbf{E}, \mathbf{H}) \quad \text{weakly-}^* \text{ in } L^\infty((0, T), L^2(\Omega) \times L^2(\Omega)) \text{ as } h \rightarrow 0, N \rightarrow \infty \\ \frac{d}{dt}(\mathbf{E}_{N,h}, \mathbf{H}_{N,h}) &\overset{*}{\rightharpoonup} \frac{d}{dt}(\mathbf{E}, \mathbf{H}) \quad \text{weakly-}^* \text{ in } L^\infty((0, T), L^2(\Omega) \times L^2(\Omega)) \text{ as } h \rightarrow 0, N \rightarrow \infty, \end{aligned}$$

where (\mathbf{E}, \mathbf{H}) is the unique solution to (P). Assume additionally that

$$\mathbf{H} \in L^1((0, T), \mathbf{H}(\text{curl})) \quad \text{and} \quad \max_{n \in \{1, \dots, N-1\}} \|\text{curl } \mathbf{H}_h^{n-\frac{1}{2}}\|_{L^2(\omega)} \leq C.$$

Then it holds that

$$(\mathbf{E}_{N,h}, \mathbf{H}_{N,h}) \rightarrow (\mathbf{E}, \mathbf{H}) \quad \text{in } C([0, T], L^2(\Omega) \times L^2(\Omega)) \text{ as } h \rightarrow 0.$$

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- Derivation of the following system for the weak limit:

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- Stability estimates \Rightarrow Existence of weakly-star converging subsequences
- Derivation of the following system for the weak limit:

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Main idea: Bypass missing stability by exploiting properties of **piecewise constant interpolation operator** for $\mathbf{v} \in \mathbf{K} \cap \mathbf{C}_0^\infty(\Omega)$.

- $(P_{\text{weak}}) \Rightarrow (P)$ reduces to **enlarging** the set of test functions

$$K \cap C_0^\infty(\Omega) \rightsquigarrow K \cap H_0(\text{curl})$$

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Question: Does there exist a mollification operator M_δ , s.t.

$$v \in K \cap H_0(\text{curl}) \stackrel{?}{\Rightarrow} \begin{cases} M_\delta v \in C_0^\infty(\Omega) \\ M_\delta v \in K. \end{cases}$$

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$$v \in K \cap H_0(\text{curl}) \stackrel{?}{\Rightarrow} \begin{cases} M_\delta v \in C_0^\infty(\Omega)^{3,4} \\ M_\delta v \in K. \end{cases}$$

³A. Ern and J.-L. Guermond. Mollification in strongly Lipschitz domains with application to continuous and discrete de Rham complexes. *Comput. Methods Appl. Math.*, 16(1):51–75, 2016

⁴S.H. Christiansen and R. Winther, Smoothed projections in finite element exterior calculus, *Math. Comp.* 77 (2008), no. 262, 813–829

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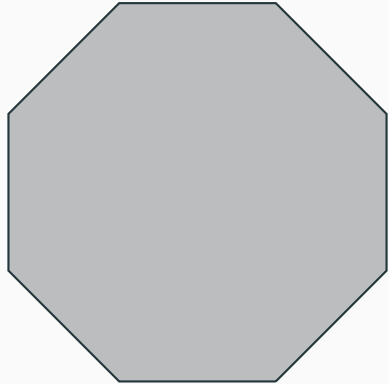
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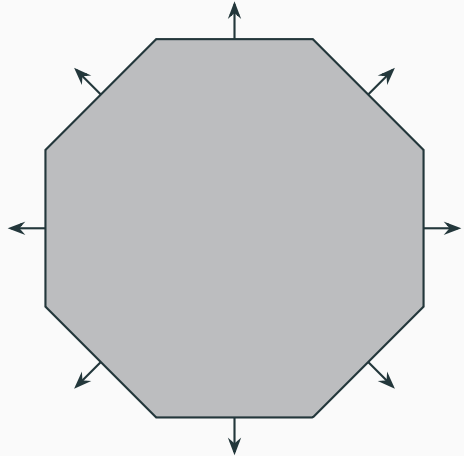
⁴S.H. Christiansen and R. Winther, Smoothed projections in finite element exterior calculus, *Math. Comp.* 77 (2008), no. 262, 813–829

Idea of Ern and Guermond: $\mathbf{v} \in H_0(\mathbf{curl})$.



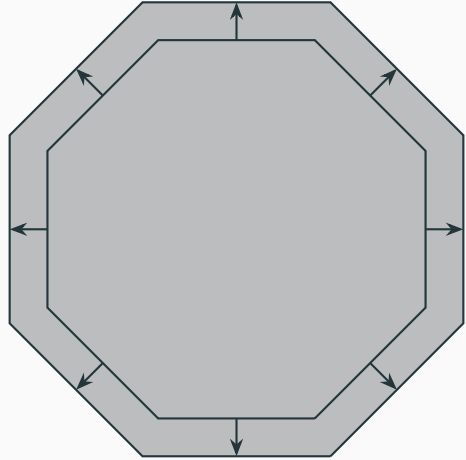
Idea of Ern and Guermond: $\mathbf{v} \in H_0(\mathbf{curl})$.

- Expand Ω by a transversal vector field



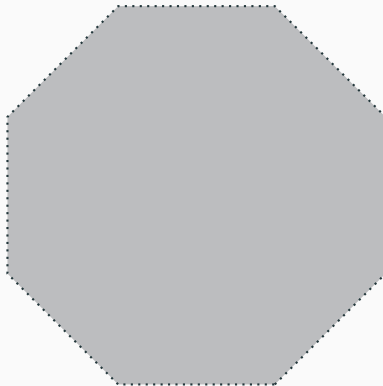
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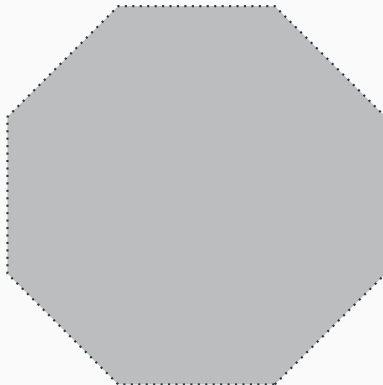
- Expand Ω by a transversal vector field
- Cut off vector field



Idea of Ern and Guermond: $\mathbf{v} \in H_0(\mathbf{curl})$.

- Expand Ω by a transversal vector field
- Cut off vector field

\rightsquigarrow Mollify resulting field $\Rightarrow \mathbf{M}_\delta \mathbf{v} \in C_0^\infty(\Omega)$

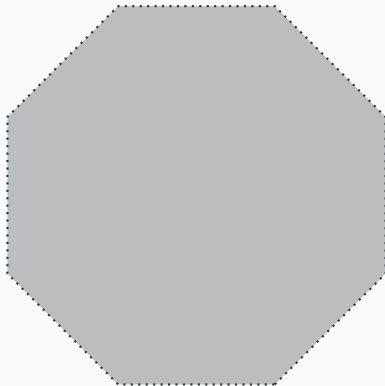


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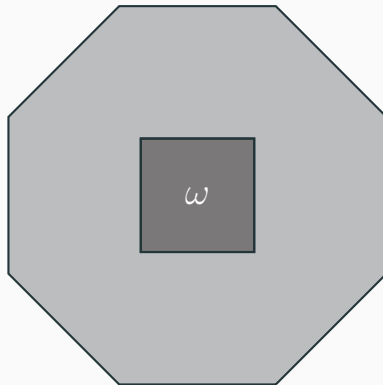
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Q: And what about points in ω ?



Idea of Ern and Guermond: $\mathbf{v} \in H_0(\mathbf{curl})$.

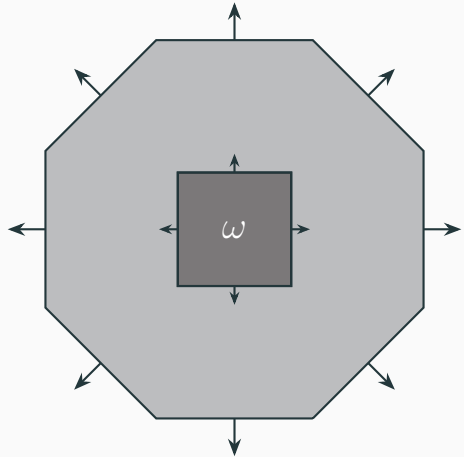
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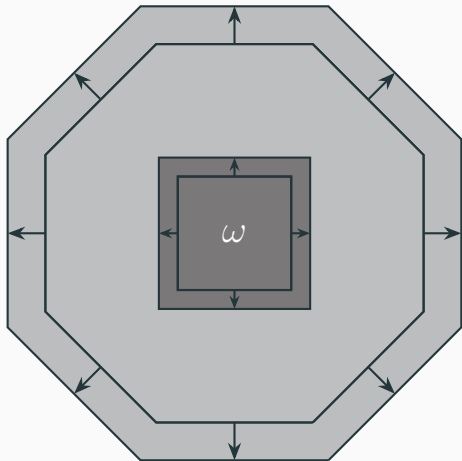
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Idea of Ern and Guermond: $\mathbf{v} \in H_0(\mathbf{curl})$.

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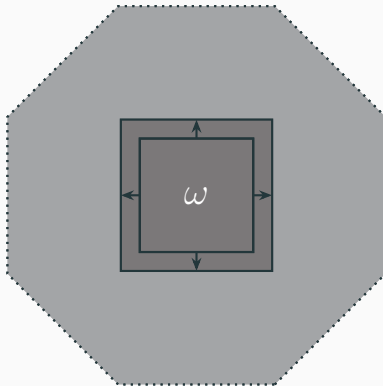
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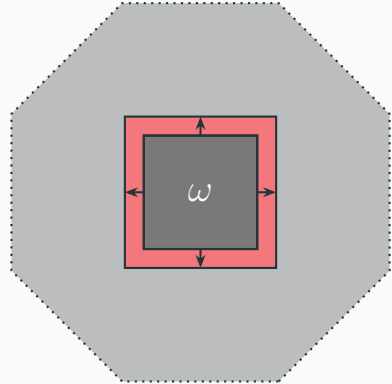


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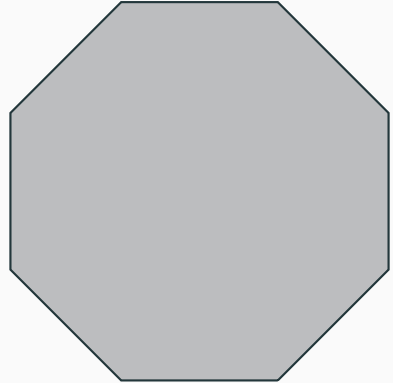
- Expand Ω by a transversal vector field
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Q: And what about points in ω ?

A: Could get pushed outside of ω .

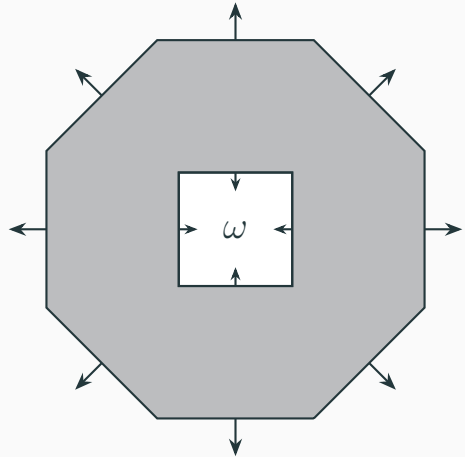


Workaround:



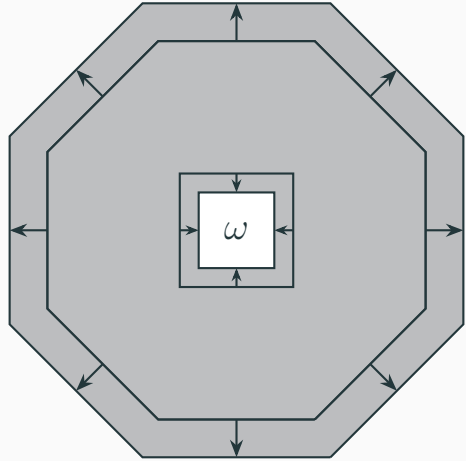
Workaround:

- This time expand $\Omega \setminus \bar{\omega}$



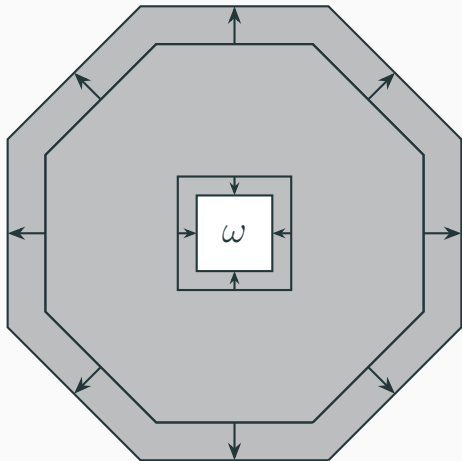
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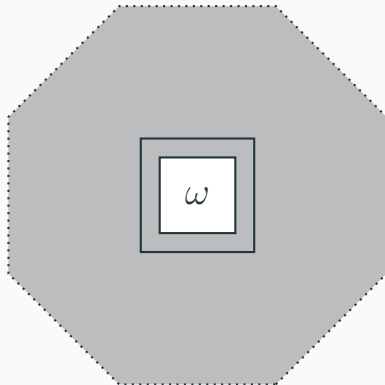
- This time expand $\Omega \setminus \bar{\omega}$
 \rightsquigarrow Possible, since $\Omega \setminus \bar{\omega}$ Lipschitz⁵



⁵S. Hofmann, M. Mitrea, and M. Taylor. Geometric and transformational properties of Lipschitz domains,... *J. Geom. Anal.*, 17(4):593–647, 2007

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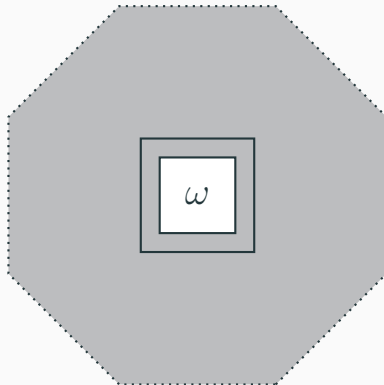
- This time expand $\Omega \setminus \bar{\omega}$
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- Cut off vector field near the boundary



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 \rightsquigarrow Mollify resulting vector field



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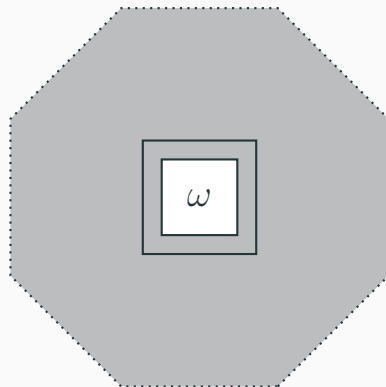
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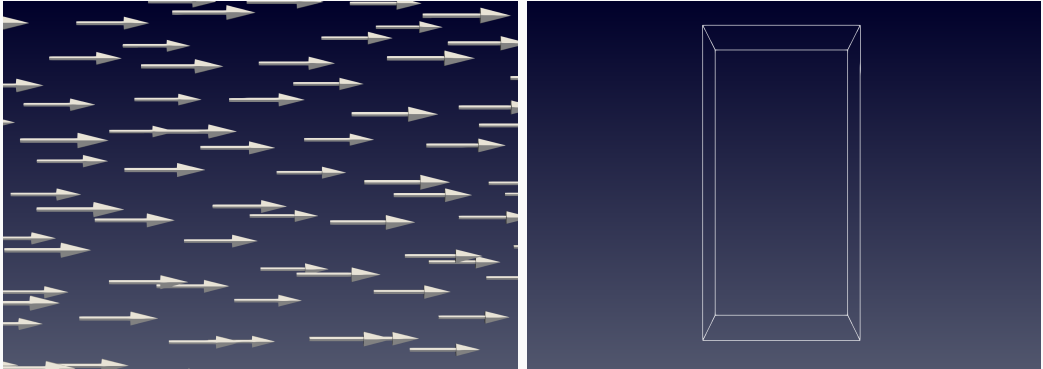
Techniques from geometrical analysis:

$$\mathbf{v} \in K \cap H_0(\mathbf{curl}) \Rightarrow M_\delta \mathbf{v} \in K$$

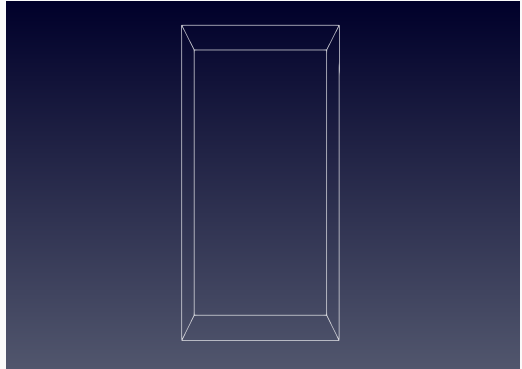
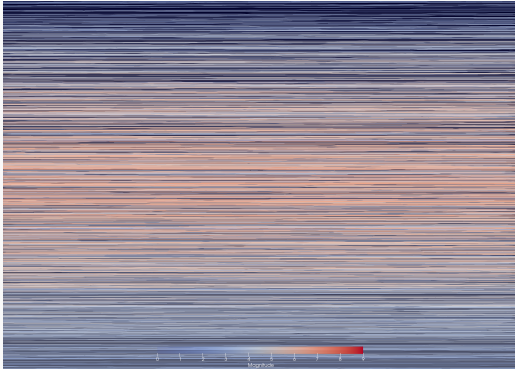


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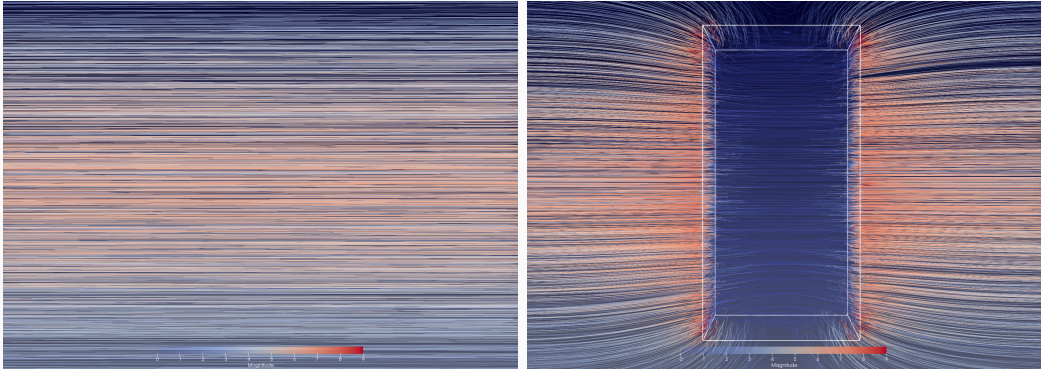
Numerical test using FEniCS



Applied current source and obstacle



Free electric field and obstacle



Free and shielded electric field

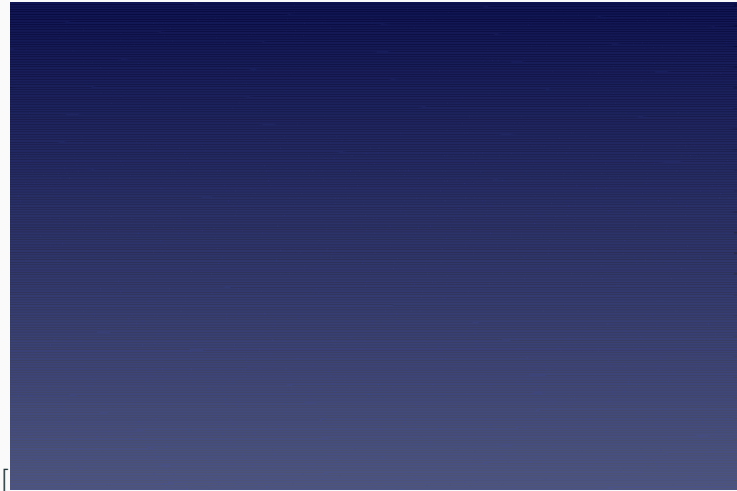
Reference parameters:

$$\tilde{h} = 1/2^6, \quad \tilde{N} = 5 \cdot 2^6 \quad \rightsquigarrow \quad \mathbf{E} := \mathbf{E}_{\tilde{N}, \tilde{h}}$$

Define:

$$\text{RelErr}_{N,h}(\mathbf{E}) := \frac{\|\mathbf{E}_{N,h} - \mathbf{E}\|_{C([0,T],L^2(\Omega))}}{\|\mathbf{E}\|_{C([0,T],L^2(\Omega))}} \approx \frac{\max_{n \in \{0, \dots, N\}} \|\mathbf{E}_{N,h}(t_n) - \mathbf{E}(t_n)\|_{L^2(\Omega)}}{\max_{n \in \{0, \dots, N\}} \|\mathbf{E}(t_n)\|_{L^2(\Omega)}}$$

N	$5 \cdot 2^2$	$5 \cdot 2^3$	$5 \cdot 2^4$	$5 \cdot 2^5$	$5 \cdot 2^6$
h	$1/2^2$	$1/2^3$	$1/2^4$	$1/2^5$	$1/2^6$
$\dim(\text{DG}_h)$	1.152	9.216	31.024	589.824	4.718.592
$\dim(\text{ND}_h)$	604	4.184	73.728	238.688	1.872.064
$\text{RelErr}_{N,h}(\mathbf{E})$	0.3832	0.1070	0.0591	0.0248	—
$\text{RelErr}_{N,h}(\mathbf{H})$	0.3556	0.2083	0.1480	0.1086	—



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Thank you for your attention!